

Part III Point-based multiphysics modeling

Daniel Holz^{1,2}, Stefan Rhys Jeske³, Fabian Löschner³, Jan Bender³, Yin Yang⁴, Sheldon Andrews²



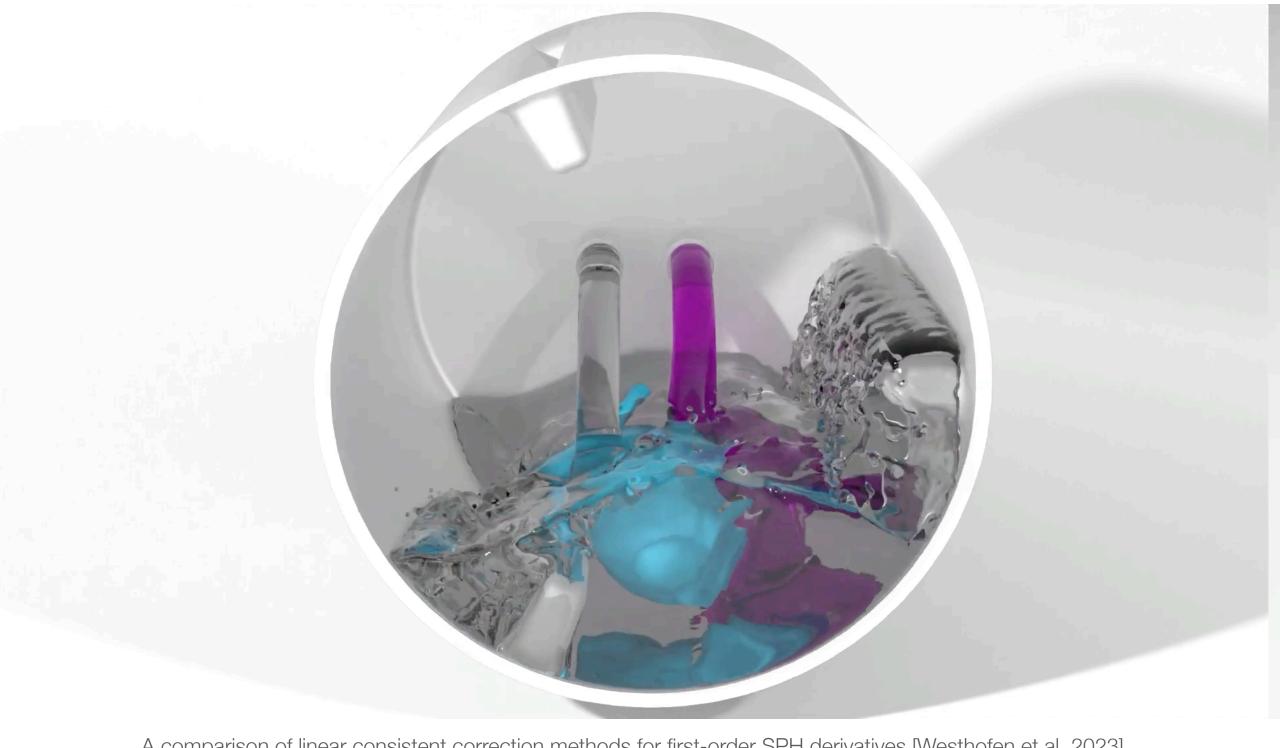






Lagrangian Point-based Methods

Eulerian and Hybrid Methods







Fast and versatile fluid-solid coupling for turbulent flow simulation [Lyu et al. 2021]





Overview

- Lagrangian Point-based Methods
 - Overview
 - Fluid
 - Solid and Rigid
 - Multiphysics Materials
- Eulerian and Hybrid Methods
 - Overview
 - Multiphysics Materials
- Summary

	Lagrangian Point-Based Methods (Sec. 2)	П	Eulerian & Hybrid Methods (Sec. 3)	Energy-Based Modeling (Sec. 4)	Constraint-Based Modeling (Sec. 5)
Deformables (elastic & plastic)	[MKN*04] [PKA*05] [SSP07] [BIT09] [MKB*10] [YJL*16] [YCL*17] [PGBT18] [CLC*20] [KBF*21] [KUKH23]		[SZS95] [CGF006] [LLJ*11] [SSJ*14] [JSS*15] [YSB*15] [TLK16] [FGG*17] [GTJS17] [JGT17] [ZB17] [GHF*18] [HFG*18] [FLGJ19] [HGG*19] [SXH*21] [LLJ22] [TB22] [QLY*23] [LLH*24] [TLZ*24]	[BAV*10] [BUAG12] [SB12b] [SHST12] [BML*14] [GSS*15] [LBK17] [BOFN18] [SGK18] [LFS*20] [MEM*20] [LMY*22] [LCK22] [LLJ22] [KE22] [LFFJ*23]	[Jak01] [MHTG05] [MHHR06] [SLM06] [MMCK14] [BKCW14] [Cho14] [MCKM15] [CMM16] [DCB16] [MMC16] [BGAO17] [FM17] [ARM*19] [MEM*19] [WWB*19] [MMC*20] [MM21] [TTKA23] [CHC*24a] [Cet24] [MAK24] [SZDJ24] [YLL*24]
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Rigid Bodies & Multibody Systems	[SSP07] [YCL*17] [GPB*19] [PT23]		[TB20] [TB22] [LLH*24] [TLZ*24]	[CDGB19] [MEM*20] [FLS*21] [CLL*22] [LKL*22]	[Bar94] [MC95] [ST96] [Bar96] [AP97] [Stc00] [Jak01] [Erl05] [MHTG05] [Lac07b, Lac07a] [GZO10] [MMCK14] [DCB16] [FM17] [MEM*19] [PAK*19] [WWB*19] [MMC*20] [MAK24]
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Fluids & Fluid Phenomena	[PW02] [MCG03] [SSP07] [BT07] [BIT09] [SP09] [Pri12] [SB12a] [AAT13] [ICS*14] [HWZ*15] [TDF*15] [BK17] [PT17] [YCL*17] [YML*17] [PGBT18] [WKBB18] [BKKW19] [CBG*19] [GPB*19] [WJL*20] [ZRS*20] [KBF*21] [LWB*21] [WDK*21] [LHWW22] [XRW*22] [JWL*23] [PT23] [XLYJ23] [ZLX*24] [YWX*24]		[Har62] [HW*65] [BR86] [FM96] [Sta99] [Pes02] [TUKF02] [CMT04] [ZB05] [CGF006] [KFC006] [CFL*07] [MCP*09] [SABS14] [SSJ*14] [ATW15] [JSS*15] [RGJ*15] [FGG*17] [GPH*18] [HFG*18] [JGT17] [ZB17] [FLGJ19] [GAB20] [HGMRT20] [TB20] [CKMR*21] [SXH*21] [QLDGJ22] [TB22] [STBA24] [QLY*23] [LLH*24] [TLZ*24]	[TB20] [TB21] [TB22] [XLYJ23]	[BLS12] [MM13] [MMCK14] [TNF14] [BGAO17] [XRW*22] [YLL*24]
Multi-Phase, Phase Transitions & Porous Flow	[MKN*04] [SSP07] [LAD08] [SP08] [BIT09] [LD09] [PC13] [RLY*14] [YCR*15] [YJL*16] [PGBT18] [CLC*20] [GHB*20] [WFM21] [RXL21] [RHLC22] [XWW*23] [YR23] [ZLX*24]		[SSJ*14] [ATW15] [GPH*18] [GAB20] [CKMR*21] [SXH*21] [LMLD22] [TLZ*24]		[MMCK14]
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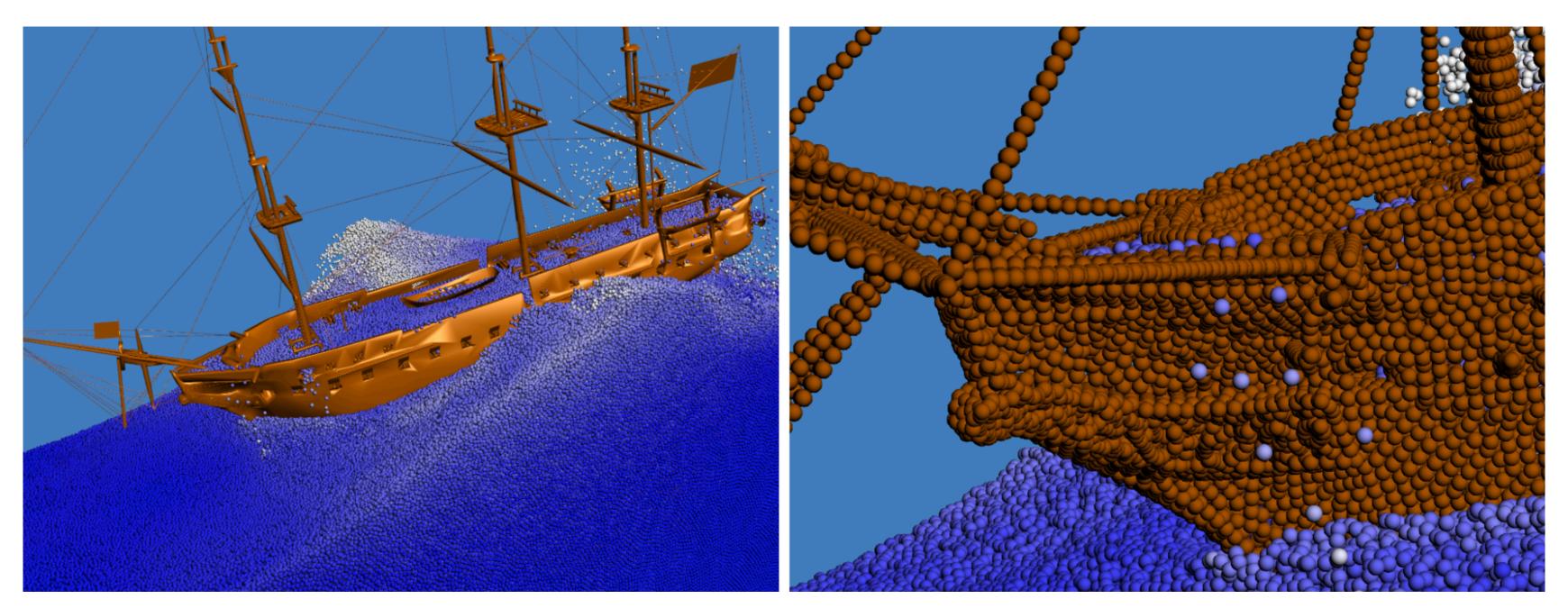




Lagrangian Point-based Methods - Overview

Simplest primitive

- Suitable for topology changes
- Versatile description of many objects
- "Locking" particle arrangements



Versatile rigid-fluid coupling for incompressible SPH [Akinci et al. 2012]







Lagrangian Point-based Methods - Method zoo

- Smoothed Particle Hydrodynamics (SPH)
- Moving-Least Squares (MLS)
- Reproducing Kernel Particle Method (RKPM)
- Discrete Element Method (DEM)*
- Moving Particle Semi-implicit Method (MPS)
- Peridynamics

- Multiphysical
- Graphics

- Multiphysical
- Graphics





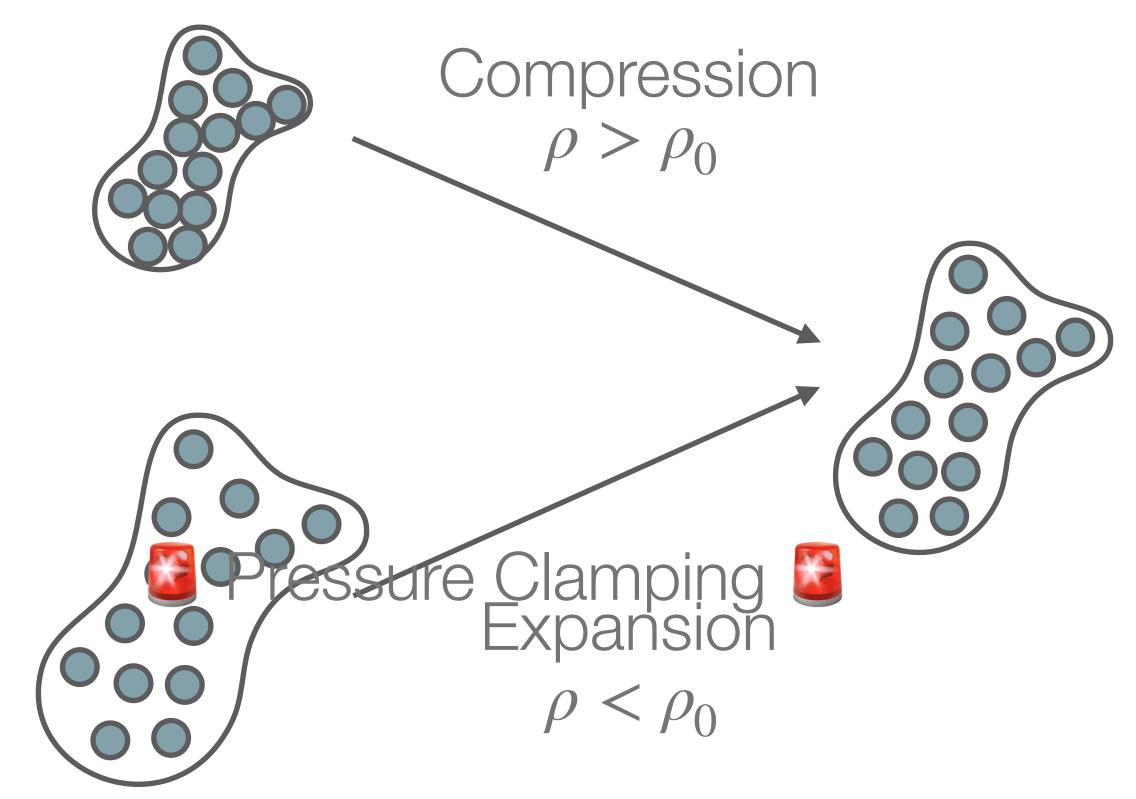
Lagrangian Point-based Methods - Equations of Motion

$$\rho \frac{D\mathbf{v}}{Dt} = -\nabla p + \mu \nabla^2 \mathbf{v} + \mathbf{f} \qquad \frac{D\rho}{Dt} = -\rho(\nabla \cdot \mathbf{v}) = 0$$

Pressure-Poisson Equation (PPE)

$$\Delta t \, \nabla^2 p = \frac{D\rho}{Dt}$$

$$\Delta t \nabla^2 p = \rho(\nabla \cdot \mathbf{v})$$









$$\rho \frac{D\mathbf{v}}{Dt} = -\nabla p + \mu \nabla^2 \mathbf{v} + \mathbf{f}$$

Pressure $-\nabla p$

• Explicit [Müller et al. 2003]

Eurographics/SIGGRAPH Symposium on Computer Animation (2003) D. Breen, M. Lin (Editors)

Particle-Based Fluid Simulation for Interactive Applications

Matthias Müller, David Charypar and Markus Gross

Department of Computer Science, Federal Institute of Technology Zürich (ETHZ), Switzerland

Abstract

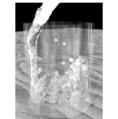
Realistically animated fluids can add substantial realism to interactive applications such as virtual surgery simulators or computer games. In this paper we propose an interactive method based on Smoothed Particle Hydrodynamics (SPH) to simulate fluids with free surfaces. The method is an extension of the SPH-based technique by Desbrun to animate highly deformable bodies. We gear the method towards fluid simulation by deriving the force density fields directly from the Navier-Stokes equation and by adding a term to model surface tension effects. In contrast to Eulerian grid-based approaches, the particle-based approach makes mass conservation equations and convection terms dispensable which reduces the complexity of the simulation. In addition, the particles can directly be used to render the surface of the fluid. We propose methods to track and visualize the free surface using point splatting and marching cubes-based surface reconstruction. Our animation method is fast enough to be used in interactive systems and to allow for user interaction with models consisting of up to 5000 particles.

Categories and Subject Descriptors (according to ACM CCS): I.3.7 [Computer Graphics]: Three-Dimensional Graphics and Realism

1. Introduction

1.1. Motivation

Fluids (i.e. liquids and gases) play an important role in every day life. Examples for fluid phenomena are wind, weather, ocean waves, waves induced by ships or simply pouring of a glass of water. As simple and ordinary these phenomena may seem, as complex and difficult it is to simulate them. Even though Computational Fluid Dynamics (CFD) is a well





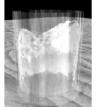


Figure 1: Pouring water into a glass at 5 frames per second.





$$\rho \frac{D\mathbf{v}}{Dt} = -\nabla p + \mu \nabla^2 \mathbf{v} + \mathbf{f}$$

Pressure $-\nabla p$

- Explicit [Müller et al. 2003]
- Constraint-based [Bodin et al. 2012] [Macklin and Müller 2013]

IEEE TRANSACTIONS OF VISUALIZATION AND COMPUTER GRAPHICS

Constraint Fluids

Kenneth Bodin, Claude Lacoursière, Martin Servin

Position Based Fluids

Miles Macklin *

Matthias Müller †

NVIDIA

Abstract

In fluid simulation, enforcing incompressibility is crucial for realism; it is also computationally expensive. Recent work has improved efficiency, but still requires time-steps that are impractical for real-time applications. In this work we present an iterative density solver integrated into the Position Based Dynamics framework (PBD). By formulating and solving a set of positional constraints that enforce constant density, our method allows similar incompressibility and convergence to modern smoothed particle hydrodynamic (SPH) solvers, but inherits the stability of the geometric, position based dynamics method, allowing large time steps suitable for real-time applications. We incorporate an artificial pressure term that improves particle distribution, creates surface tension, and lowers the neighborhood requirements of traditional SPH. Finally, we address the issue of energy loss by applying vorticity confinement as a velocity post process.

CR Categories: I.3.5 [Computer Graphics]: Computational Geometry and Object Modeling—Physically based modeling I.3.7 [Computer Graphics]: Three-Dimensional Graphics and Realism—

Keywords: fluid simulation, SPH, PCISPH, constraint fluids, position based dynamics

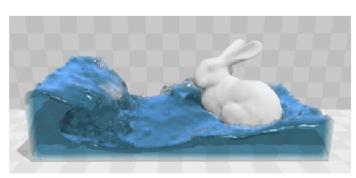
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1 Introduction

Fluids, in particular liquids such as water, are responsible for many visually rich phenomena, and simulating them has been an area of long-standing interest and challenge in computer graphics. There are a variety of techniques available, but here we focus on particle methods, which are popular for their simplicity and flexibility.

Smoothed Particle Hydrodynamics (SPH) [Monaghan 1992][1994] is a well known particle based method for fluid simulation. It



(a) Real-time rendered fluid surface using ellipsoid splatting



(b) Underlying simulation particles

Figure 1: Bunny taking a bath. 128k particles, 2 sub-steps, 3 density iterations per frame, average simulation time per frame 10ms.

jarola 2009], but small time steps remain a requirement, limiting real-time applications.

For interactive environments, robustness is a key issue: the simulation must handle degenerate situations gracefully. SPH algorithms often become unstable if particles do not have enough neighbors for accurate density estimates. The typical solution is to try to avoid these situations by taking sufficiently small time steps, or by using sufficiently many particles, at the cost of increased computation.

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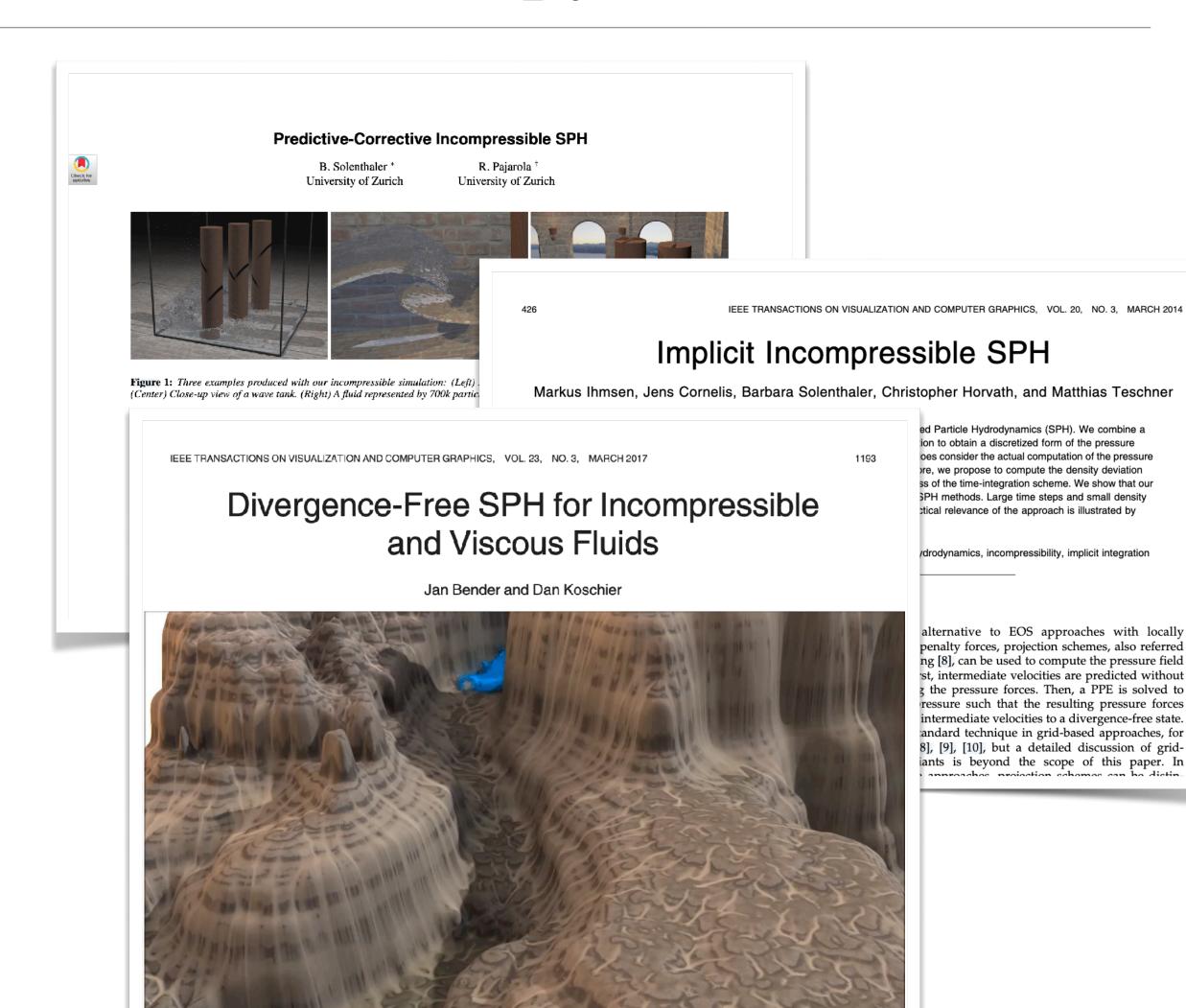




$$\rho \frac{D\mathbf{v}}{Dt} = -\nabla p + \mu \nabla^2 \mathbf{v} + \mathbf{f}$$

Pressure $-\nabla p$

- Explicit [Müller et al. 2003]
- Constraint-based [Bodin et al. 2012] [Macklin and Müller 2013]
- Implicit [Solenthaler and Pajarola 2009] [Ihmsen et al. 2014] [Bender et al. 2017]





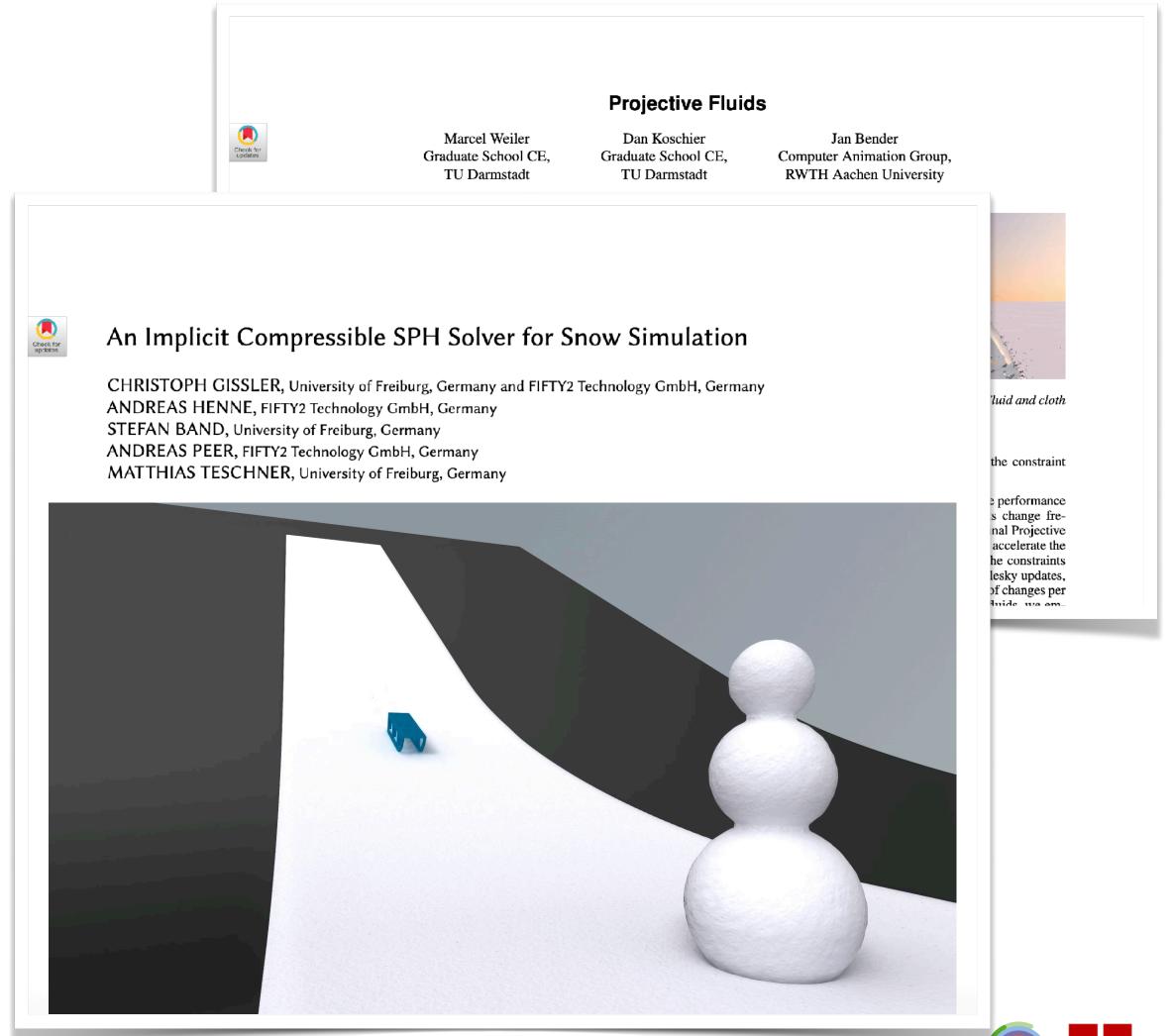




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- Compressible [Gissler et al. 2020] [Weiler et al. 2016]









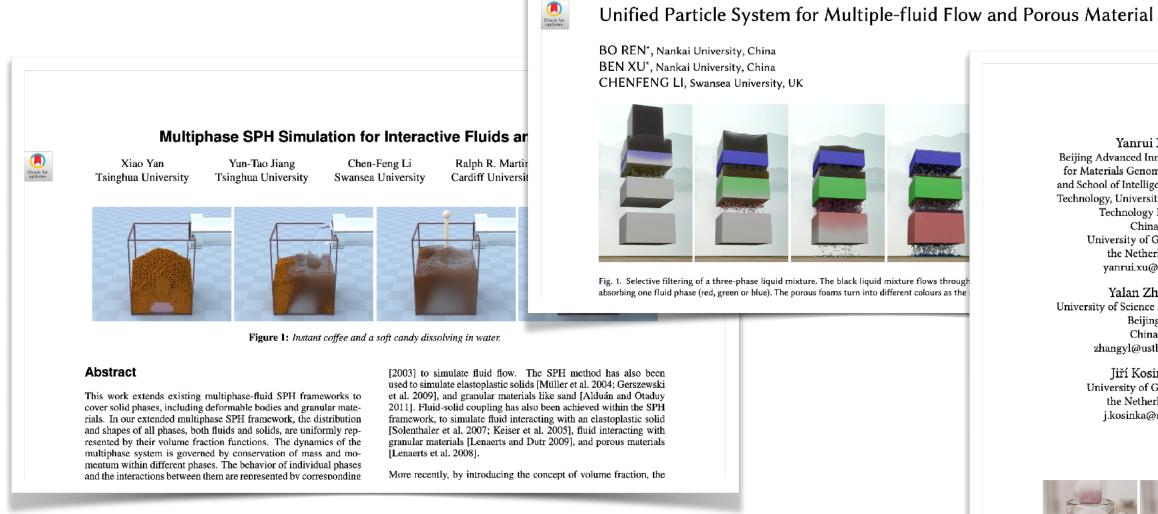
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Density Contrast SPH Interfaces

B. Solenthaler and R. Pajarola Visualization and MultiMedia Lab, University of Zurich, Switzerland

- Multiphase
 - Multi-fluid [Solenthaler and Pajarola 2008]

• Mixing and dissolution [Xu et al. 2023] [Ren et al. 2021] [Yan et al. 2016]



An Implicitly Stable Mixture Model for Dynamic Multi-fluid Simulations

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$$\rho \frac{D\mathbf{v}}{Dt} = -\nabla p + \mu \nabla^2 \mathbf{v} + \mathbf{f}$$

- Multiphase
 - Multi-fluid [Solenthaler and Pajarola 2008]
 - Mixing and dissolution [Xu et al. 2023] [Ren et al. 2021] [Yan et al. 2016]
- Turbulence
 - Micropolar [Bender et al. 2019]
 - Vorticity Refinement Monte Carlo [Ye et al. 2024] [Liu et al. 2021]





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Turbulent Micropolar SPH Fluids with Foam

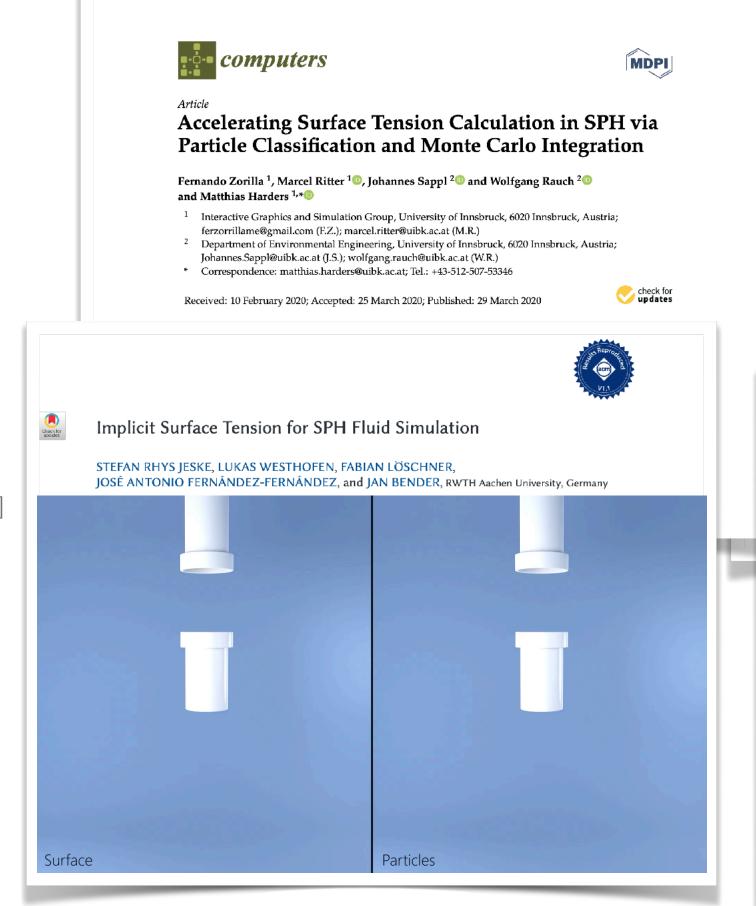
Jan Bender[®], Dan Koschier[®], Tassilo Kugelstadt, and Marcel Weiler



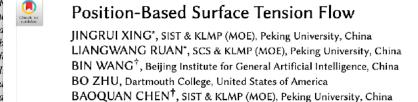


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 - Vorticity Refinement Monte Carlo [Ye et al. 2024] [Liu et al. 2021]
- Surface Tension
 - Curvature [Xing et al. 2022] [Zorilla et al. 2020] [Müller et al. 2003]
 - Cohesion [Jeske et al. 2023] [Yang et al. 2017]







Eurographics/SIGGRAPH Symposium on Computer Animation (2003)



Pairwise Force SPH Model for Real-Time Multi-Interaction Applications

Tao Yang, Ralph R. Martin, Ming C. Lin, Fellow, IEEE, Jian Chang, and Shi-Min Hu

Abstract—In this paper, we present a novel pairwise-force smoothed particle hydrodynamics (PF-SPH) model to enable simulation of various interactions at interfaces in real time. Realistic capture of interactions at interfaces is a challenging problem for SPH-based simulations, especially for scenarios involving multiple interactions at different interfaces. Our PF-SPH model can readily handle multiple types of interactions simultaneously in a single simulation; its basis is to use a larger support radius than that used in standard SPH. We adopt a novel anisotropic filtering term to further improve the performance of interaction forces. The proposed model is stable; furthermore, it avoids the particle clustering problem which commonly occurs at the free surface. We show how our model can be used to capture various interactions. We also consider the close connection between droplets and bubbles, and show how to animate bubbles rising in liquid as well as bubbles in air. Our method is versatile, physically plausible and easy-to-implement. Examples are provided to demonstrate the capabilities and effectiveness of our approach

Index Terms—Smoothed particle hydrodynamics (SPH), pairwise force, surface tension, bubble animation, fluid simulation

1 INTRODUCTION

TN computer graphics, interactions at interfaces between high density ratios [2] and interfacial flows [3], [4], [6]. We erials in different phases, or immiscible materials in avoid complex treatments for miscible flows and conc the same phase, have been extensively investigated during on the interactions between different immiscible flows. The the last decade. In this work, we focus on interfaces involving interactions between a liquid and a solid are of two main a liquid; thus commonly observed interfaces can be categorized into three classes: between a liquid and a gas, another pling contributes to macroscopic movements while adhesion liquid, or a solid. The gas is mostly considered as air in this is caused by molecular forces and results in various wetting paper unless otherwise specified. The interaction between a effects. In the real world, many scenarios involve interactions liquid and air leads to surface tension, which is the main at multiple interfaces. For instance, when cracking an egg, cause of many well known visual effects, including the water there are simultaneous interactions between air, egg white.







- Sample with Particles
 - Force-based [Becker et al. 2009]
 - Pressure-based [Akinci et al. 2013] [Akinci et al. 2012]
 - Rigid-rigid & rigid-fluid [Probst and Teschner 2023] [Gissler et al. 2019]



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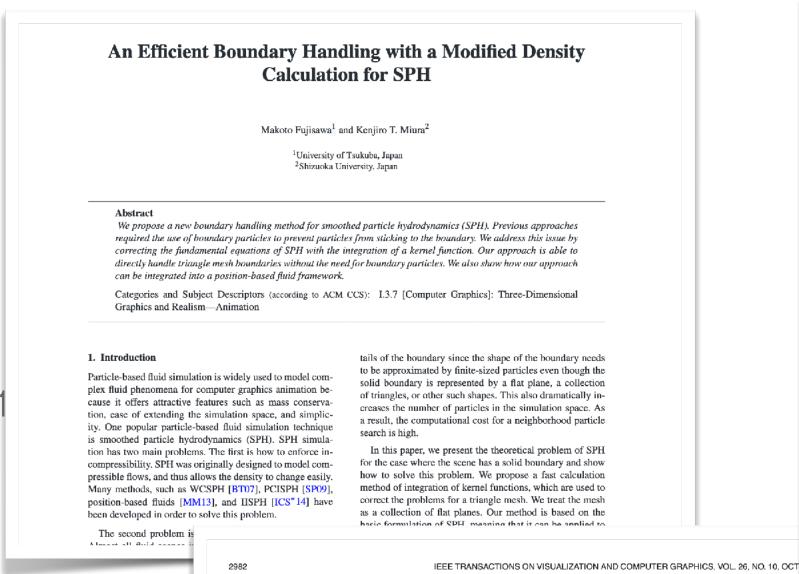






applications and simulation econories. Henally in combined an

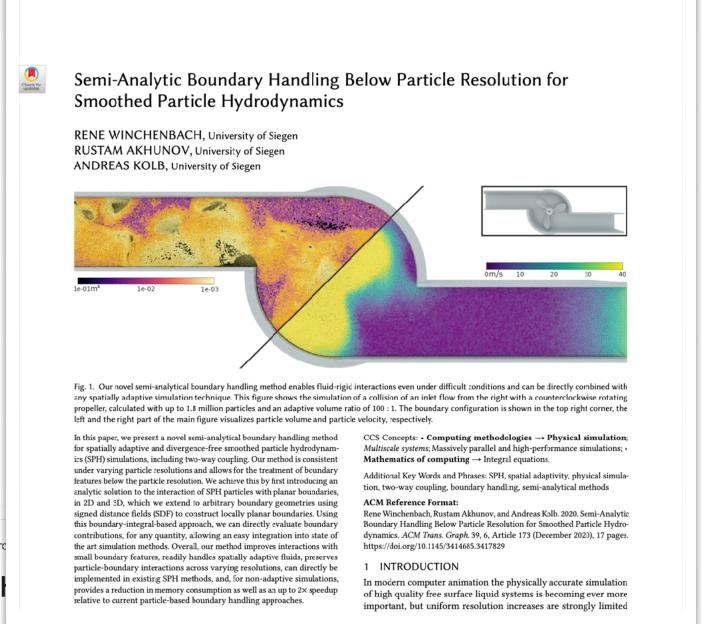
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 - Pressure-based [Akinci et al. 201 [Akinci et al. 2012]
 - Rigid-rigid & rigid-fluid [Probst and Teschner 2023] [Gissler et al. 2019]
- Implicit boundaries [Winchenbach et al. 2020] [Bender et al. 2020] [Fujisawa and Miura 2015]



Implicit Frictional Boundary Handling for SP

Jan Bender[®], Tassilo Kugelstadt[®], Marcel Weiler[®], and Dan Koschier[®]

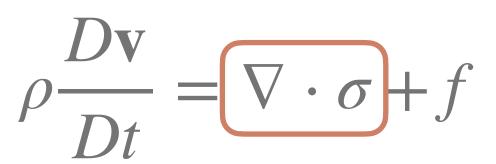












- Elasticity
 - Linear [Solenthaler et al. 2007]
 - Corotational [Becker et al. 2009]
 - Implicit coronated linear [Kugelstadt et al. 2021] [Peer et al. 2018]
 - Non-linear [Kee et al. 2023]

COMPUTER ANIMATION AND VIRTUAL WORLDS Comv. Anim. Virtual Worlds 2007: 18: 69-82 Published online in Wiley InterScience (www.interscience.wiley.com) DOI: 10.1002/cav.162 A unified particle model for fluid-soli Eurographics Workshop on Natural Phenomena (2009) interactions E. Galin and J. Schneider (Editors) By Barbara Solenthaler*, Jürg Schläfli and Renato Pajarola We present a new method for the simulation of melting and solidification in a unified particle model. Our technique uses the Smoothed Particle Hydrodynamics (SPH) meth the simulation of liquids, deformable as well as rigid objects, which eliminates the ne define an interface for coupling different models. Using this approach, it is possible to simulate fluids and solids by only changing the attribute values of the underlying par-We significantly changed a prior elastic particle model to achieve a flexible model for melting and solidification. By using an SPH approach and considering a new definitio local reference shape, the simulation of merging and splitting of different objects, as m Smoothed Particle Hydrodynamics (SPH) is a powerful techn caused by phase change processes, is made possible. In order to keep the system stable early SPH approaches in Computer Graphics have mainly b in regions represented by a sparse set of particles we use a special kernel function for also focuses on the dynamics of deformable solids using solidification processes. Additionally, we propose a surface reconstruction technique l SPH formulation for deformable solids. The rigid body m on considering the movement of the center of mass to reduce rendering errors in concar allows to use a linear strain tensor. In contrast to previous examples using conlanar and collinear particle data sets. EUROGRAPHICS 2023 / K. Myszkowski and M. Nießner COMPUTER GRAPHICS forum Volume 42 (2023), Number 2 An Optimization-based SPH Solver for Simulation of Hyperelastic Solids Min Hyung Kee¹, Kiwon Um², HyunMo Kang¹ and JungHyun Han¹ ¹Korea University, Korea ²LTCI, Telecom Paris, IP Paris, France

Figure 1: Our optimization-based solver allows stable and robust simulations of coupling between hyperelastic solids and fluids in a unified

SPH framework. The elastic solids of this example are simulated using the Neo-Hookean model with 39.3K particles, and the fluid is simulated

This paper proposes a novel method for simulating hyperelastic solids with Smoothed Particle Hydrodynamics (SPH). The

materials, such as the Neo-Hookean and the St. Venant-Kirchoff models. To this end, we reformulate an implicit integration

scheme for SPH elastic solids into an optimization problem and solve the problem using a general-purpose quasi-Newton

method. Our experiments show that the Limited-memory BFGS (L-BFGS) algorithm can be employed to efficiently solve our optimization problem in the SPH framework and demonstrate its stable and efficient simulations for complex materials in the

SPH framework. Thanks to the nature of our unified representation for both solids and fluids, the SPH formulation simplifies

proposed method extends the coverage of the state-of-the-art elastic SPH solid method to include different types of hyperelastic

using the divergence-free SPH solver with 320K particles.

coupling between different materials and handling collisions.

Computing methodologies → Physical simulation, elastic body simulation, optimization.

Corotated SPH for d

Markus Becker Markus Ihms

olids and fluids. This and the efficiency of

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An Implicit SPH Formulation for Incompressible Linearly

Andreas Peer (D), Christoph Gissler, Stefan Band and Matthias Teschner

Elastic Solids

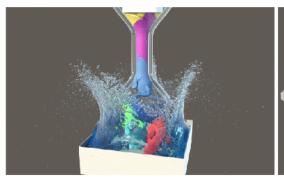
University of Freiburg, Germany {peera, gisslere, bands, teschner}@informatik.uni-freiburg.de

DOI: 10.1111/cgf.13317

We propose a novel smoothed particle hydrodynamics (SPH) formulation for deformable solids. Key aspects of our method are implicit elastic forces and an adapted SPH formulation for the deformation gradient that—in contrast to previous work—allows

Fast Corotated Elastic SPH Solids with Implicit Zero-Energy **Mode Control**

TASSILO KUGELSTADT, RWTH Aachen University, Germany JAN BENDER, RWTH Aachen University, Germany JOSÉ ANTONIO FERNÁNDEZ-FERNÁNDEZ, RWTH Aachen University, Germany STEFAN RHYS JESKE, RWTH Aachen University, Germany FABIAN LOSCHNER, RWTH Aachen University, Germany ANDREAS LONGVA, RWTH Aachen University, Germany



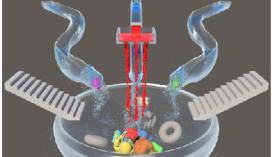


Fig. 1. Left: Stable simulation of eight walrus models (210k particles) that are pushed through a tight funnel and impact the water in a container with 1.2M fluid particles. Right: To showcase the coupling capabilities of our method 10 deformable solids and 4 rigid tori are dropped into a bowl while water and a highly viscous fluid are poured on top. A total of 800k particles are used for the fluids and 252k for the elastic objects.

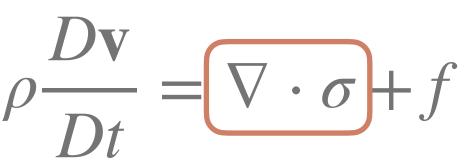
We develop a new operator splitting formulation for the simulation of corotated linearly elastic solids with Smoothed Particle Hydrodynamics (SPH). Based on the technique of Kugelstadt et al. [2018] originally developed for the Finite Element Method (FEM), we split the elastic energy into two separate terms corresponding to stretching and volume conservation, and based on this principle, we design a splitting scheme compatible with SPH. The operator splitting scheme enables us to treat the two terms separately, and because the stretching forces lead to a stiffness matrix that is constant in time, we are able to prefactor the system matrix for the implicit integration step. Solid-solid contact and fluid-solid interaction is achieved through a unified pressure solve. We demonstrate more than an order of magnitude improvement in computation time compared to a





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COMPUTER GRAPHICS forum



- Elasticity
 - Linear [Solenthaler et al. 2007]
 - Corotational [Becker et al. 2009]
 - Implicit coronated linear [Kugelstadt et al. 2021] [Peer et al. 2018]
 - Non-linear [Kee et al. 2023]
- Plasticity
 - Mohr Coulomb [Lenaerts and Dutré 2009]
 - Drucker Prager [Alduán and Otaduy 2011] [Ihmsen et al. 2013]
 - Snow [Gissler et al. 2020]









Overview

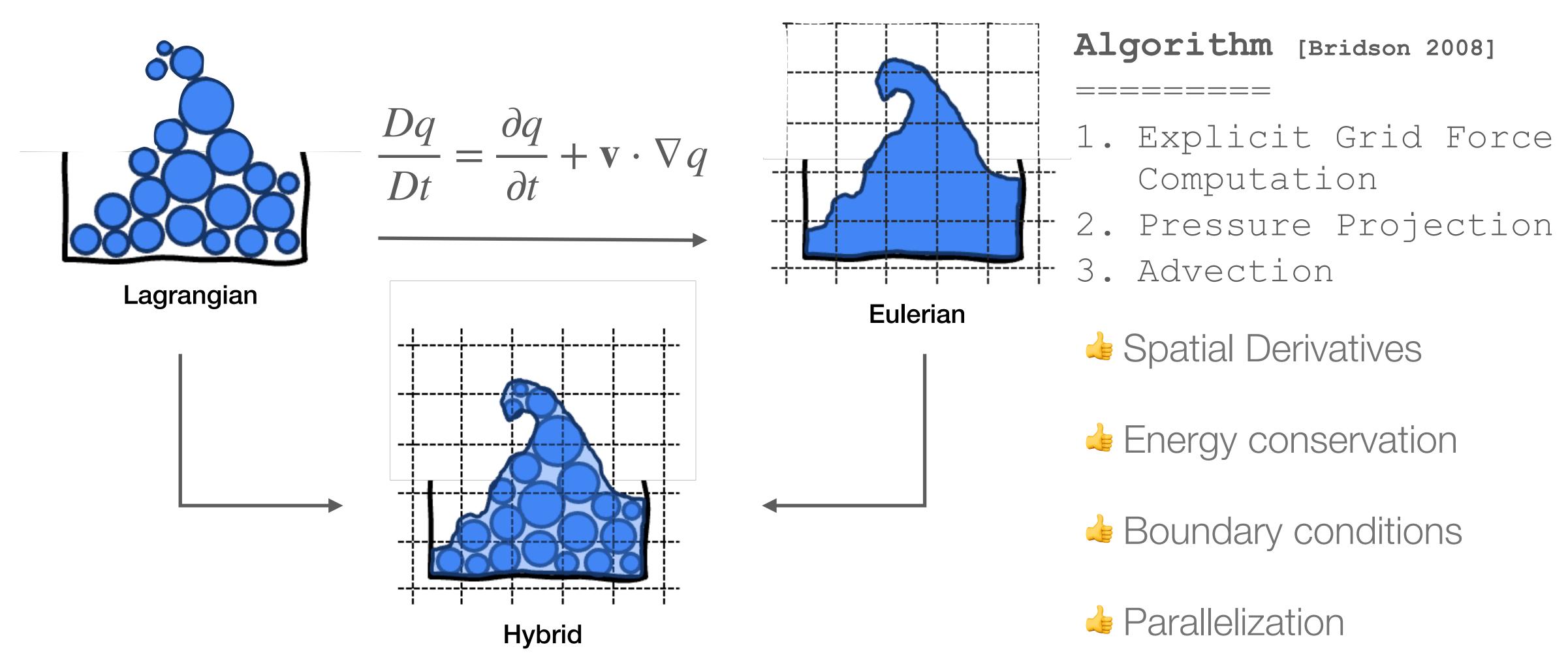
- Lagrangian Point-based Methods
 - Overview
 - Fluid
 - Solid and Rigid
 - Multiphysics Materials
- Eulerian and Hybrid Methods
 - Overview
 - Multiphysics Materials
- Summary

	Lagrangian Point-Based Methods (Sec. 2)	Eulerian & Hybrid Methods (Sec. 3)	Energy-Based Modeling (Sec. 4)	Constraint-Based Modeling (Sec. 5)
Deformables (elastic & plastic)	[MKN*04] [PKA*05] [SSP07] [BIT09] [MKB*10] [YJL*16] [YCL*17] [PGBT18] [CLC*20] [KBF*21] [KUKH23]	[SZS95] [CGFO06] [LLJ*11] [SSJ*14] [JSS*15] [YSB*15] [TLK16] [FGG*17] [GTJS17] [JGT17] [ZB17] [GHF*18] [HFG*18] [FLGJ19] [HGG*19] [SXH*21] [LLJ22] [TB22] [QLY*23] [LLH*24] [TLZ*24]	[BAV*10] [BUAG12] [SB12b] [SHST12] [BML*14] [GSS*15] [LBK17] [BOFN18] [SGK18] [LFS*20] [MEM*20] [LMY*22] [LCK22] [LLJ22] [KE22] [LFFJ*23]	[Jak01] [MHTG05] [MHHR06] [SLM06] [MMCK14] [BKCW14] [Cho14] [MCKM15] [CMM16] [DCB16] [MMC16] [BGAO17] [FM17] [ARM*19] [MEM*19] [WWB*19] [MMC*20] [MM21] [TTKA23] [CHC*24a] [Cet24] [MAK24] [SZDJ24]
Granular Materials	[LD09] [AO11] [IWT13] [YJL*16] [YCL*17] [GHB*20]	[ZB05] [SSC*13] [DBD16] [KGP*16] [TGK*17] [GPH*18]		[Hol14] [MMCK14] [SWLB14] [FM17] [HG18] [NS18] [KKHS20] [YLL*24]
Rigid Bodies & Multibody Systems	[SSP07] [YCL*17] [GPB*19] [PT23]	[TB20] [TB22] [LLH*24] [TLZ*24]	[CDGB19] [MEM*20] [FLS*21] [CLL*22] [LKL*22]	[Bar94] [MC95] [ST96] [Bar96] [AP97] [Stc00] [Jak01] [Erl05] [MHTG05] [Lac07b, Lac07a] [GZO10] [MMCK14] [DCB16] [FM17] [MEM*19] [PAK*19] [WWB*19] [MMC*20] [MAK24]
Co-dimensional Structures	[MKB*10] [ZQC*14] [ZLQF15]	[JGT17] [GHF*18] [HGG*19] [LLH*24]	[GHDS03] [ST07] [BWR*08] [CSvRV18] [Kim20] [LKJ21] [CXY*23] [HB23] [SWP*23] [WB23] [LFFJB24]	[Jak01] [MHHR06] [GHF*07] [SL08] [SLNB10] [MKC12] [BKCW14] [MMCK14] [USS15] [MMC16] [KS16] [DKWB18] [ARM*19]
Fluids & Fluid Phenomena	[PW02] [MCG03] [SSP07] [BT07] [BIT09] [SP09] [Pri12] [SB12a] [AAT13] [ICS*14] [HWZ*15] [TDF*15] [BK17] [PT17] [YCL*17] [YML*17] [PGBT18] [WKBB18] [BKKW19] [CBG*19] [GPB*19] [WJL*20] [ZRS*20] [KBF*21] [LWB*21] [WDK*21] [LHWW22] [XRW*22] [JWL*23] [PT23] [XLYJ23] [ZLX*24] [YWX*24]	[Har62] [HW*65] [BR86] [FM96] [Sta99] [Pes02] [TUKF02] [CMT04] [ZB05] [CGF006] [KFC006] [CFL*07] [MCP*09] [SABS14] [SSJ*14] [ATW15] [JSS*15] [RGJ*15] [FGG*17] [GPH*18] [HFG*18] [JGT17] [ZB17] [FLGJ19] [GAB20] [HGMRT20] [TB20] [CKMR*21] [SXH*21] [QLDGJ22] [TB22] [STBA24] [QLY*23] [LLH*24] [TLZ*24]	[TB20] [TB21] [TB22] [XLYJ23]	[BLS12] [MM13] [MMCK14] [TNF14] [BGAO17] [XRW*22] [YLL*24]
Multi-Phase, Phase Transitions & Porous Flow	[MKN*04] [SSP07] [LAD08] [SP08] [BIT09] [LD09] [PC13] [RLY*14] [YCR*15] [YJL*16] [PGBT18] [CLC*20] [GHB*20] [WFM21] [RXL21] [RHLC22] [XWW*23] [YR23] [ZLX*24]	[SSJ*14] [ATW15] [GPH*18] [GAB20] [CKMR*21] [SXH*21] [LMLD22] [TLZ*24]		[MMCK14]
Other Phenomena	[LL10] [Pri12]	[WFL*19] [WDG*19] [SNZ*21] [FCK22] [CCL*22]	[CSvRV18] [CNZ*22] [WFFJB24]	[GZO10] [Cho14] [BCK*22]





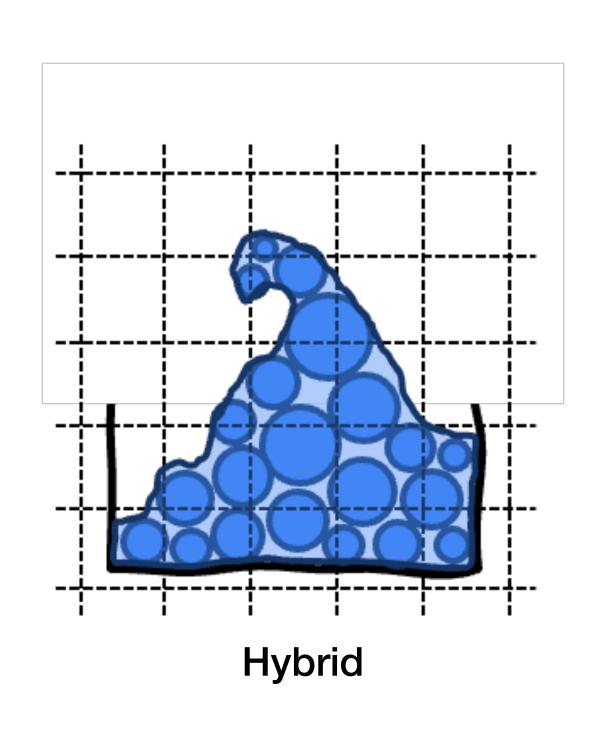
Eulerian and Hybrid Methods - Overview







Eulerian and Hybrid Methods - Overview



Lagrangian © Eulerian = MPM

- Mass conservation
- Transient grid
- Parallelization
- GPU suitable

Algorithm [Jiang et al. 2016]

- 1. Particle-to-Grid
- 2. Grid Velocity
 Computation
- 3. Explicit Grid Force Computation
- 4. Grid Velocity Update
- 5. Deformation Gradient Update
- 6. Grid-to-Particle
 Transfer
- 7. Particle Advection





• Rasterize [Takahashi et al. 2002] [Foster and Metaxas 1996]

Realistic Animation of Liquids

Nick Foster and Dimitri Metaxas

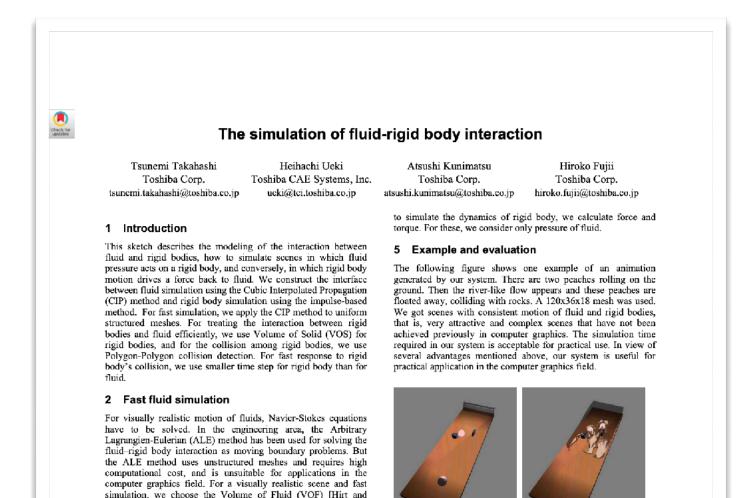
Center for Human Modeling and Simulation, University of Fennsylvania, Philadelphia, PA 19104 fostern@graphics.cis.upenn.edu http://www.cis.upenn.edu/~fostern/liquids.html

We present a comprehensive methodology for realistically animating liquid phenomena. Physically accurate 3D mo tion is achieved by performing a two-stage calculation over an arbitrary environment of static obstacles surrounded by fluid. A finite difference approximation to the Navier-Stokes equations is first applied to a low resolution, voxelized representation of the scene. The resulting velocity and pressure fields describe the gross transport of liquid, including effects such as splashing, vorticity and overturning. Local fluid velocity is then used to drive a height field equation or to convect massless marker particles. The position of any free surface can thus be determined to a significantly higher resolution than that of the Navier-Stokes calculation. In addition, the pressure field, together with the Lagrange equations of motion, is used to simulate dynamic buoyant objects. Typical disadvantages to volumetric methods such as poor scalability and lack of control are addressed by assuming that stationary obstacles align with grid cells during the finite difference discretization, and by appending driving functions to the Navier-Stokes equations. The output from our system is suitable for many of generally involve direct simulation techniques to get accuthe water rendering algorithms presented by researchers rate fluid motion. Unfortunately, in any direct simulation

Kennords: Fluid Simulations. Nanier-Stokes Equations.

effects responsible for the much of a fluid's characteristic behavior. They also cannot easily incorporate dynamic objects or buoyant effects into the model, because the ve locity of the fluid is known only on the surface, and in ternal pressure is not calculated at all. Chen and Lobo go further towards a physics-based fluid methodology by solving a simplified form of the Navier-Stokes equation in two dimensions 1. However, they assume that the fluid has zero depth, and calculate the elevation of the surface solely from the instantaneous pressure. This allows them to perform some interaction between moving objects and the flow field, but restricts the class of problems that can must be two-dimensional, and although the surface heigh is varied for animation, they treat the fluid as being completely flat during the calculation. Therefore, convective not covered by their technique.

Comprehensive models of fluid motion do exist, and technique the temporal resolution is strongly coupled to the spatial resolution. Thus, if the spatial resolution dou-

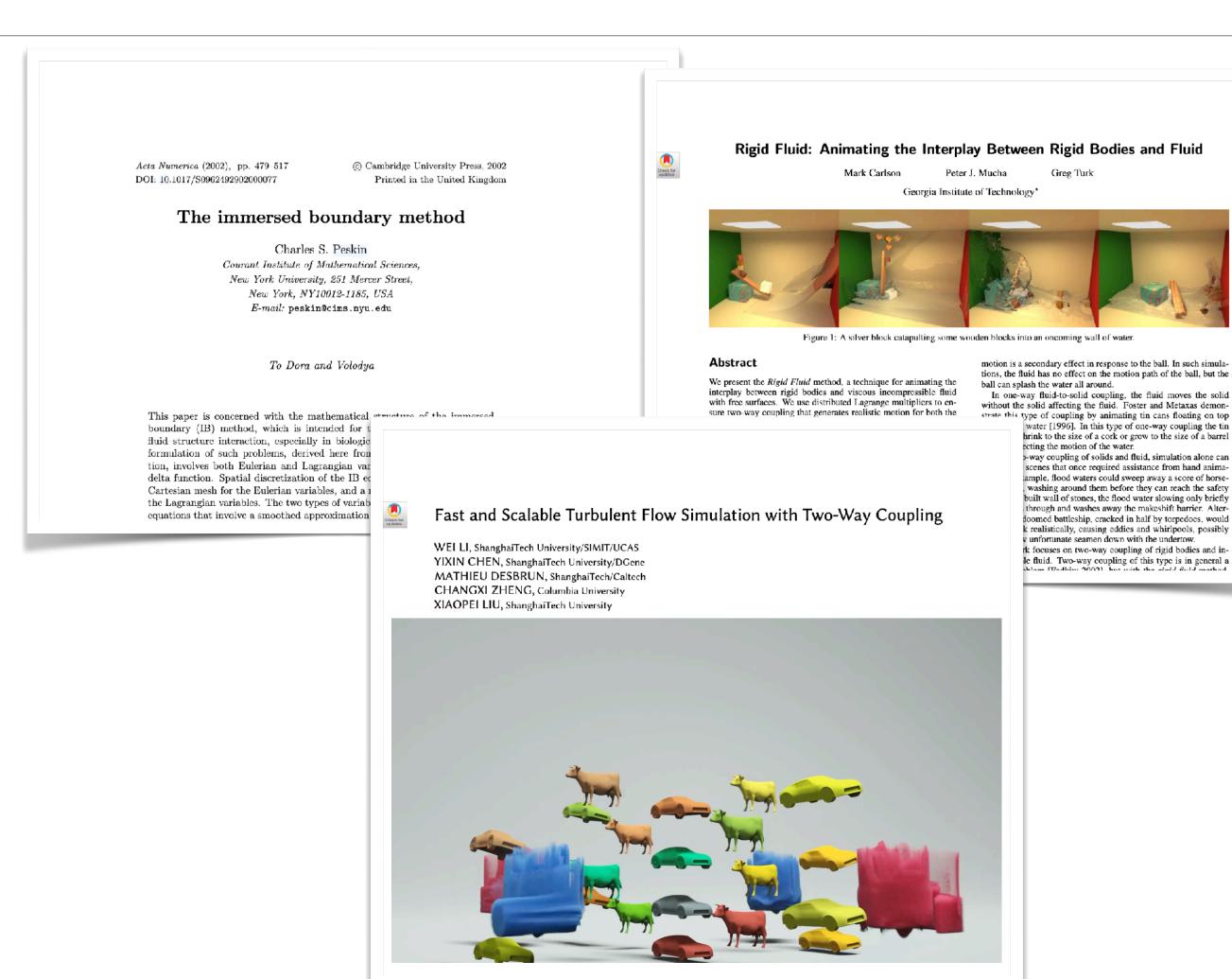


Nichols 1981] method and the CIP [Yabe 1997] method for calculation of the advection term and apply these methods to





- Rasterize [Takahashi et al. 2002] [Foster and Metaxas 1996]
- Immersed boundary method [Li et al. 2020] [Carlson et al. 2004] [Peskin 2002]

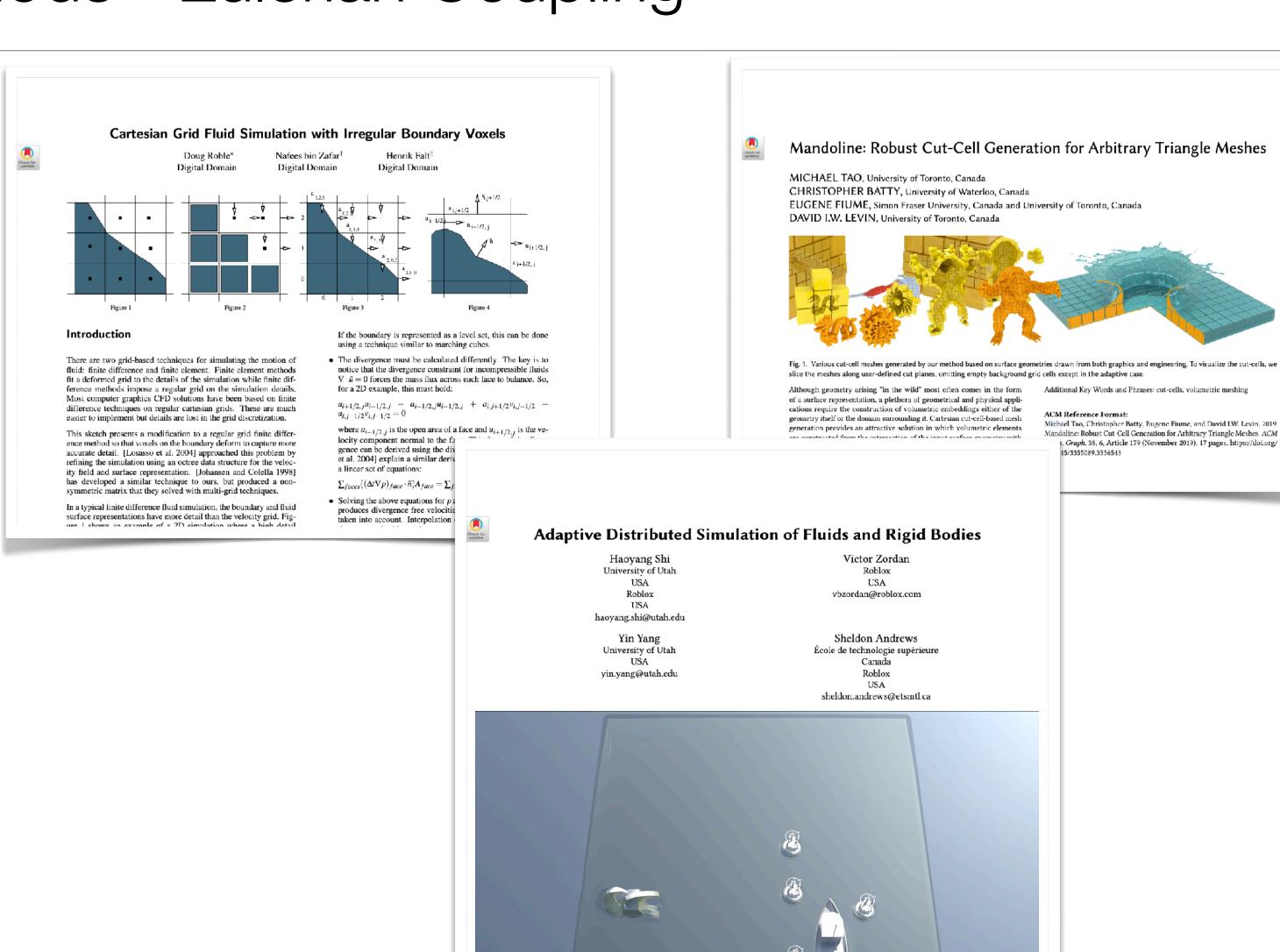








- Rasterize [Takahashi et al. 2002] [Foster and Metaxas 1996]
- Immersed boundary method [Li et al. 2020] [Carlson et al. 2004] [Peskin 2002]
- Cut-Cell [Shi et al. 2024] [Tao et al. 2019] [Roble et al. 2005]



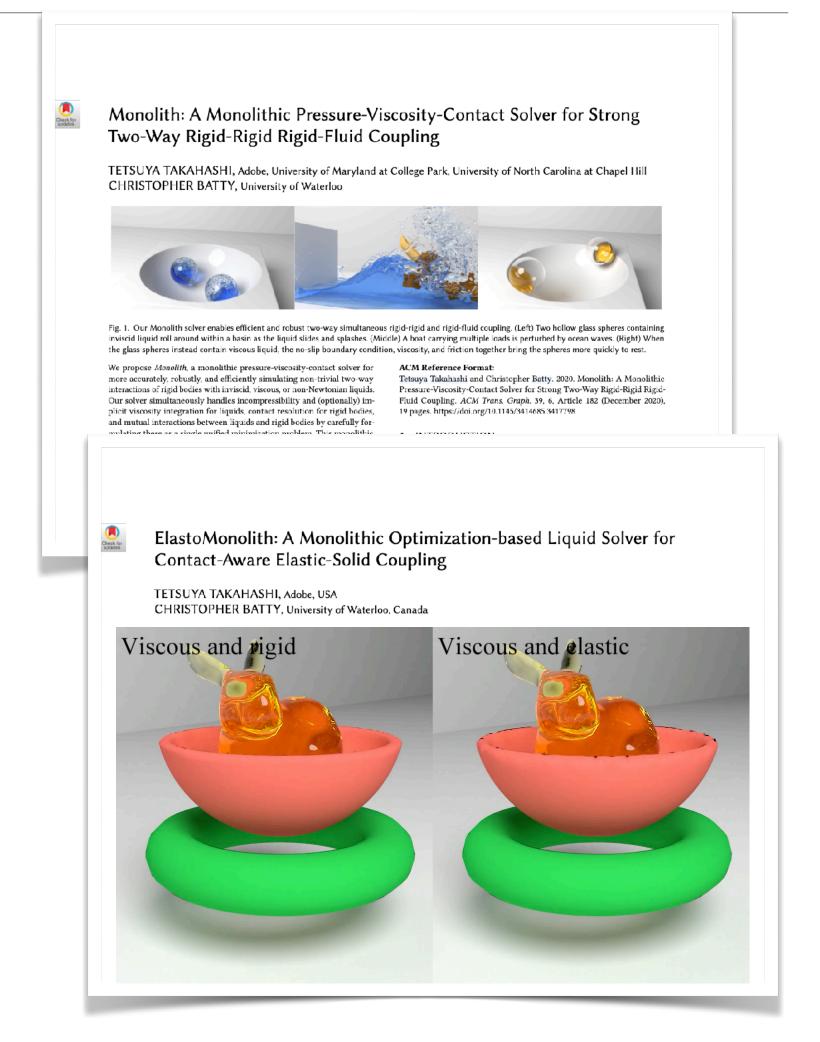






- Rasterize [Takahashi et al. 2002] [Foster and Metaxas 1996]
- Immersed boundary method [Li et al. 2020] [Carlson et al. 2004] [Peskin 2002]
- Cut-Cell [Shi et al. 2024] [Tao et al. 2019] [Roble et al. 2005]
- Rigid-fluid coupling
 - Partitioned [Akbay et al. 2018]
 - Monolithic [Takahashi and Batty 2022] [Takahashi and Batty 2020]





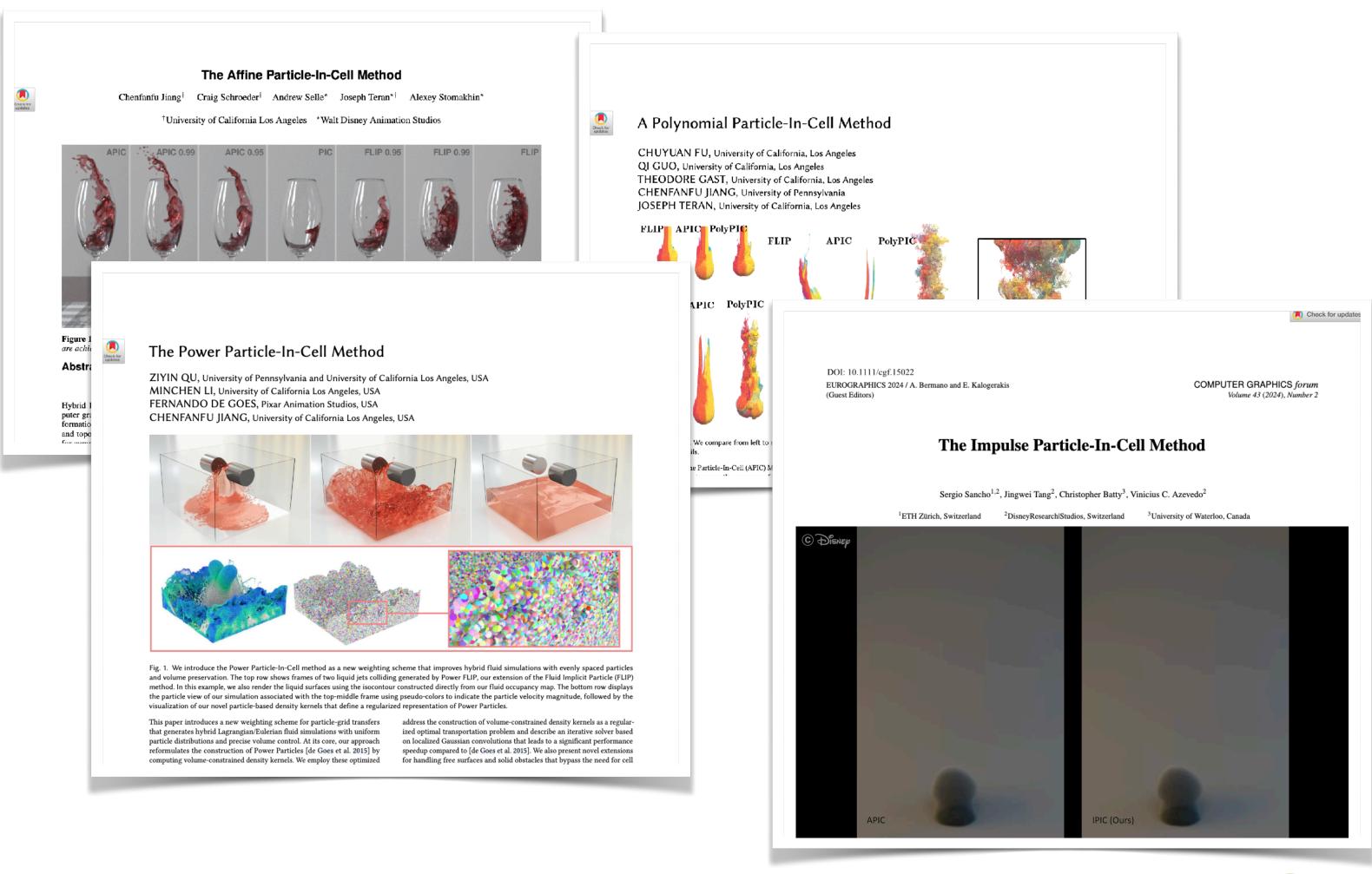






Eulerian and Hybrid Methods - Hybrid Stability and Multiphysics

- Point-to-grid transfers
 - Affine PIC [Jiang et al. 2015]
 - Poly PIC [Fu et al. 2017]
 - Power PIC [Qu et al. 2022]
 - Impulse PIC [Sancho et al. 2024]



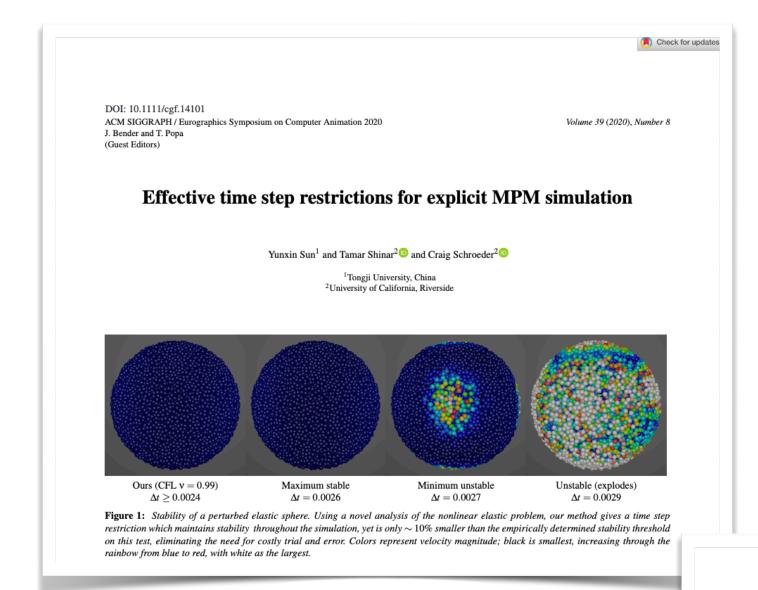






Eulerian and Hybrid Methods - Hybrid Stability and Multiphysics

- Point-to-grid transfers
 - Affine PIC [Jiang et al. 2015]
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 - Power PIC [Qu et al. 2022]
 - Impulse PIC [Sancho et al. 2024]
- Time-stepping
 - Explicit [Bai and Schroeder 2022] [Wang et al. 2020] [Sun et al. 2020]



Hierarchical Optimization Time Integration for CFL-Rate MPM Stepping XINLEI WANG, Zhejiang University & University of Pennsylvania MINCHEN LI, University of Pennsylvania & Adobe Research YU FANG, University of Pennsylvania XINXIN ZHANG, Tencent MING GAO, Tencent & University of Pennsylvania MIN TANG, Zhejiang University DANNY M. KAUFMAN, Adobe Research CHENFANFU JIANG, University of Pennsylvania We propose Hierarchical Optimization Time Integration (HOT) for effi-ACM Reference format: Xinlei Wang, Minchen Li, Yu Fang, Xinxin Zhang, Ming Gao, Min Tang, cient implicit timestepping of the material point method (MPM) irrespec Danny M. Kaufman, and Chenfanfu Jiang. 2020. Hierarchical Optimization tive of simulated materials and conditions. HOT is an MPM-specialized hierarchical optimization algorithm that solves nonlinear timestep prob-Time Integration for CFL-Rate MPM Stepping. ACM Trans. Graph. 39, 3, lems for large-scale MPM systems near the CFL limit. HOT provides con-Article 21 (April 2020), 16 pages vergent simulations out of the box across widely varying materials and computational resolutions without parameter tuning. As an implicit MPM INTRODUCTION in a quasi-Newton solver, HOT is both highly parallelizable and robustly

DOI: 10.1111/cgf.14620 ACM SIGGRAPH / Eurographics Symposium on Computer Animation 2022 D. L. Michels and S. Pirk

(Guest Editors)

Volume 41 (2022), Number 8

The material point method (MPM) is a versatile and highly effec-

tive approach for simulating widely varying material behaviors

ranging from stiff elastodynamics to viscous flows (e.g., see Fig-

ures 12 and 14) in a common framework. As such, MPM offers

the promise of a single unified, consistent, and predictive solver

for simulating continuum dynamics across diverse and potentially

heterogenous materials. However, to reach this promise, signifi

cant hurdles remain. Most significantly, obtaining accurate, consis-

Stability analysis of explicit MPM

convergent. As we show in our analysis, HOT maintains consistent and ef-

ficient performance even as we grow stiffness, increase deformation, and

vary materials over a wide range of finite strain, elastodynamic, and plastic $\,$

fectiveness of HOT against seemingly plausible alternative combinations of

MPM with standard multigrid and other Newton-Krylov models. We show

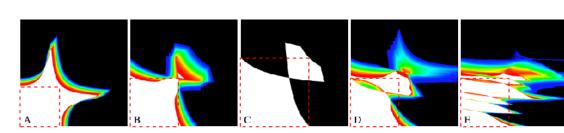
how these alternative designs result in severe issues and poor performance

In contrast, HOT outperforms existing state-of-the-art, heavily optimized

examples. Through careful benchmark ablation studies, we compare the ef-

Song Bai¹ and Craig Schroeder^{1†}

¹University of California, Riverside



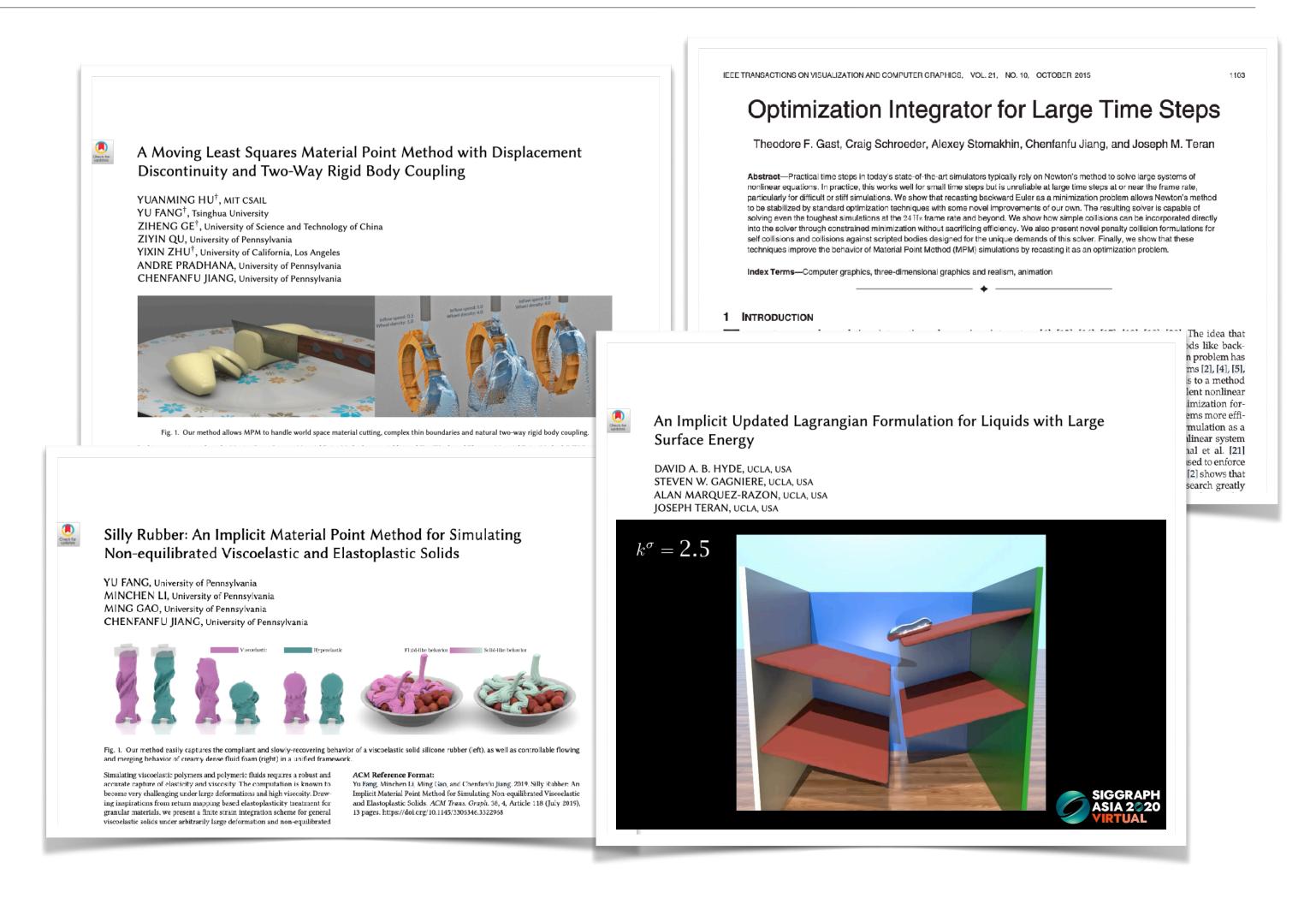
dashed box indicates the stable region based on constant time steps sizes. For images A-C, time steps alternate: $\Delta t_0, \Delta t_1, \Delta t_0, \Delta t_1, \ldots$ with PIC (A), APIC (B), and CPIC (C). In (D), APIC is run with $\Delta t_0, \Delta t_1, \Delta t_0, \Delta t_1, \Delta t_0, \Delta t_1, \Delta t_0$ is followed by 6 time steps at Δt_1 using APIC. In all cases, quadratic splines are used in 3D. All plots are on the same scale. White is stable, and black is unstable. If stability is particle position dependent, colors indicate the likelihood of any particular position being stable, with red being likely stable and blue being likely unstable. The stability region can become quite complex with unstable time step sequences scattered throughout the predicted stable region.





Eulerian and Hybrid Methods - Hybrid Stability and Multiphysics

- Point-to-grid transfers
 - Affine PIC [Jiang et al. 2015]
 - Poly PIC [Fu et al. 2017]
 - Power PIC [Qu et al. 2022]
 - Impulse PIC [Sancho et al. 2024]
- Time-stepping
 - Explicit [Bai and Schroeder 2022] [Wang et al. 2020] [Sun et al. 2020]
 - Implicit [Hyde et al. 2020] [Fang et al. 2019] [Hu et al. 2018] [Gast et al. 2015]









Eulerian and Hybrid Methods - Hybrid Materials

• Snow [Stomakhin et al. 2013]

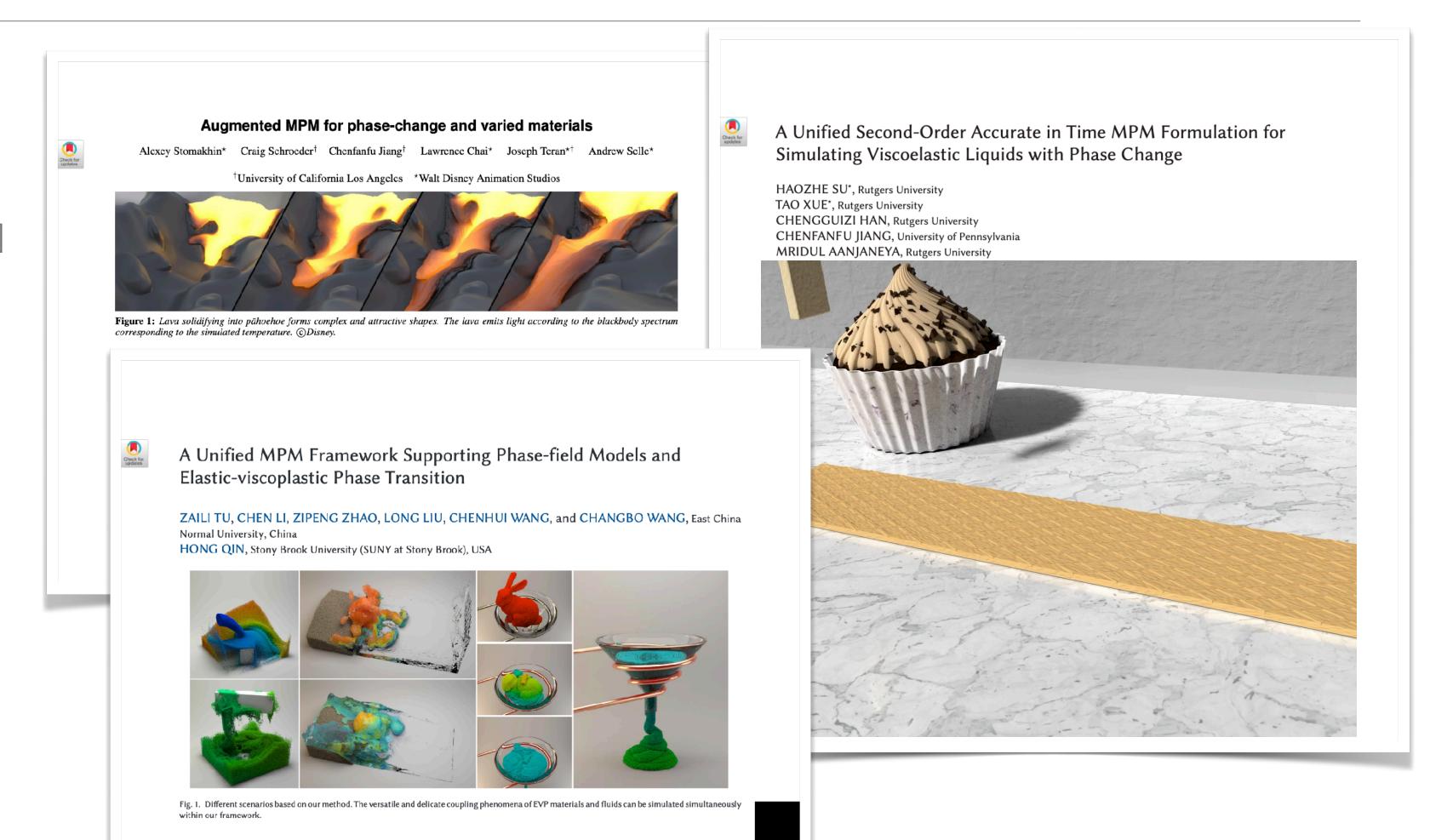






Eulerian and Hybrid Methods - Hybrid Materials

- Snow [Stomakhin et al. 2013]
- Phase-change [Tu et al. 2024] [Su et al. 2021] [Stomakhin et al. 2014]



interactions inside the same framework. In this article, we propose a prac-

tical method capable of simulating granular flows, viscoplastic liquids,

elastic-plastic solids, rigid bodies, and interacting with each other, to sup-

port novel phenomena all heavily involving realistic phase transitions,

including dissolution, melting, cooling, expansion, shrinking, and so on.







Recent years have witnessed the rapid deployment of numerous physics-

based modeling and simulation algorithms and techniques for fluids, solids,

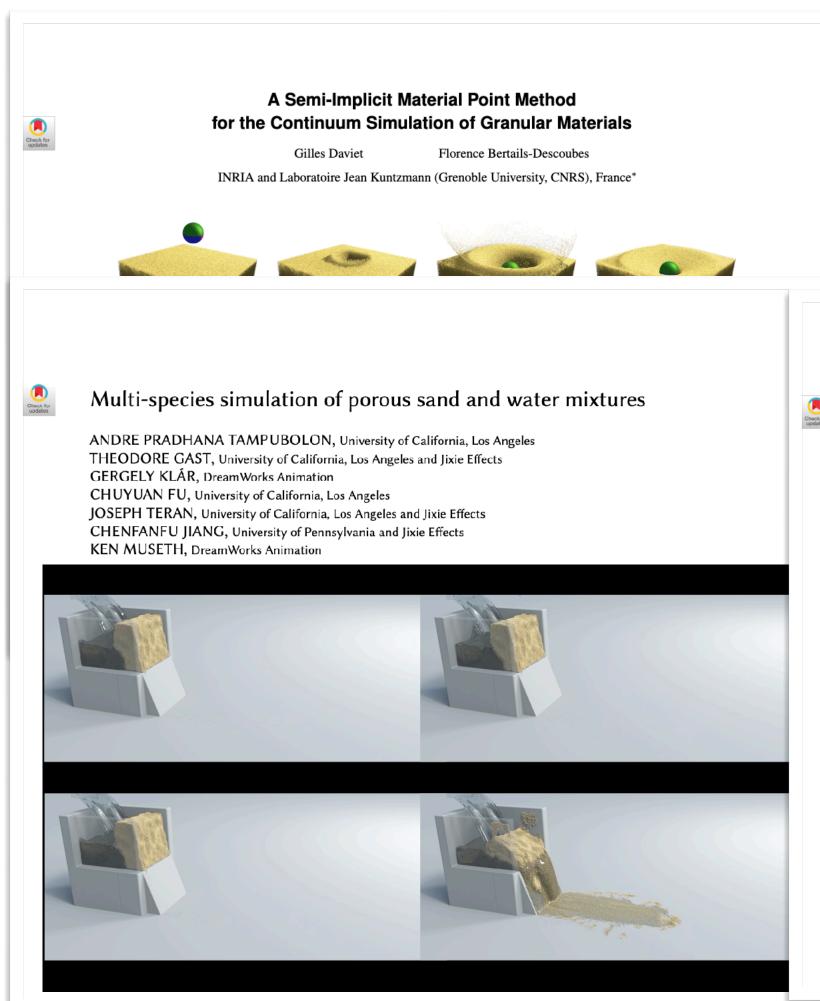
and their delicate coupling in computer animation. However, it still re-

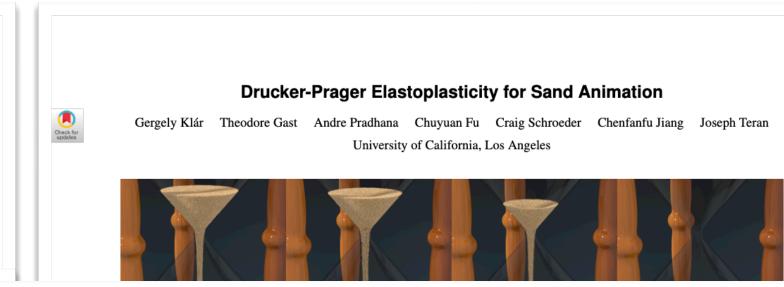
mains a challenging problem to model the complex elastic-viscoplastic

behaviors during fluid-solid phase transitions and facilitate their seamless

Eulerian and Hybrid Methods - Hybrid Materials

- Snow [Stomakhin et al. 2013]
- Phase-change [Tu et al. 2024] [Su et al. 2021] [Stomakhin et al. 2014]
- Granular
 [Li et al. 2022] [Tampubolon et al. 2017]
 [Klár et al. 2016]
 [Daviet and Bertails-Descoubes 2016]





Energetically Consistent Inelasticity for Optimization Time Integration

XUAN LI, University of California, Los Angeles, USA MINCHEN LI, University of California, Los Angeles & TimeStep Inc., USA CHENFANFU JIANG, University of California, Los Angeles & TimeStep Inc., USA

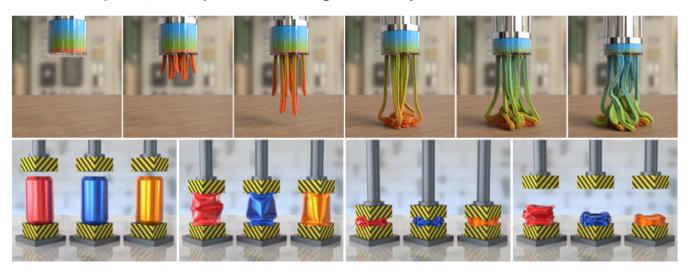


Fig. 1. Our energetically consistent inelasticity model can not only be applied to the Material Point Method (top row), but also easily extend to the Finite Element Method (bottom row with decreasing hardening coefficients from left to right). The stability under large time steps is guaranteed by the optimization time integration

In this paper, we propose Energetically Consistent Inelasticity (ECI), a new formulation for modeling and discretizing finite strain elastoplasticity/viscoelasticity in a way that is compatible with optimization-based time integrators. We provide an in-depth analysis for allowing plasticity to be implicitly integrated through an augmented strain energy density function. We develop ECI on the associative von-Mises J2 plasticity, the non-associative Drucker-Prager plasticity, and the finite strain viscoelasticity. We demonstrate the resulting scheme on both the Finite Element Method (FEM) and the Material Point Method (MPM). Combined with a custom Newton-type optimization integration scheme, our method enables simulating stiff and large-deformation inelastic dynamics of metal, sand, snow, and foam with larger time steps, improved stability, higher efficiency, and better accuracy than existing approaches.

ACM Reference Format

Xuan Li, Minchen Li, and Chenfanfu Jiang. 2022. Energetically Consistent Inelasticity for Optimization Time Integration. ACM Trans. Graph. 41, 4, Article 52 (July 2022), 16 pages. https://doi.org/10.1145/3528223.3530072

1 INTRODUCTION

Since the pioneering work of Terzopoulos and Fleischer [1988], the computer graphics community has observed increasing interests in modeling inelastic deformations governed by elastoplasticity, viscoelasticity, and viscoplasticity. These inelastic mechanical properties govern the behaviors of a wide range of everyday objects.

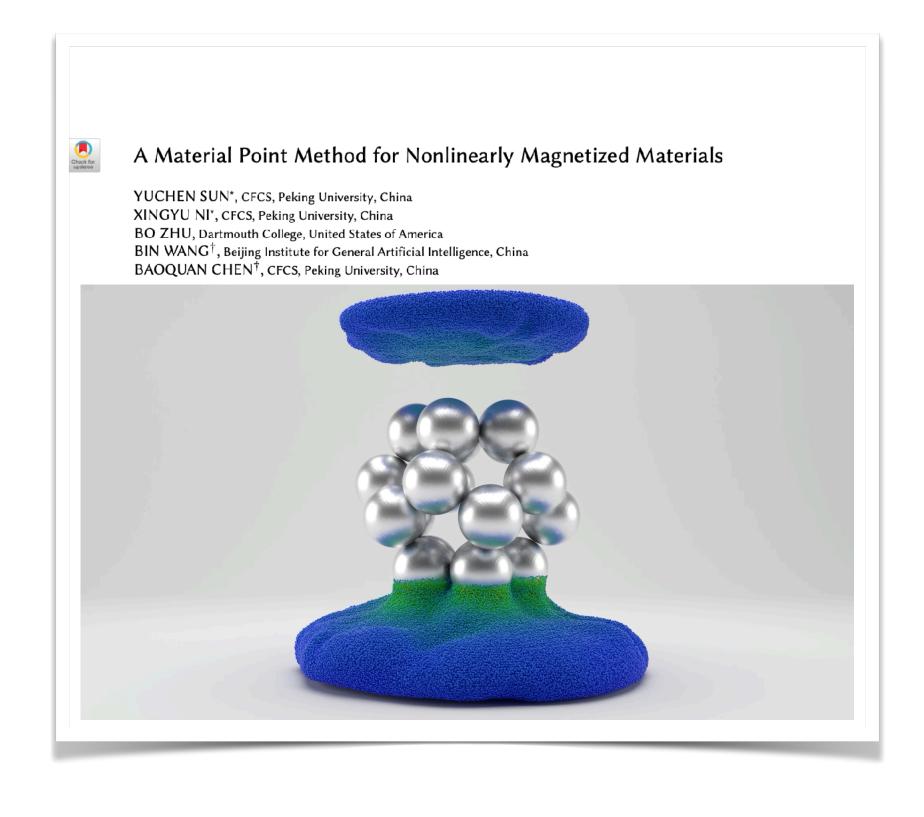






Eulerian and Hybrid Methods - Hybrid Multiphysics Materials

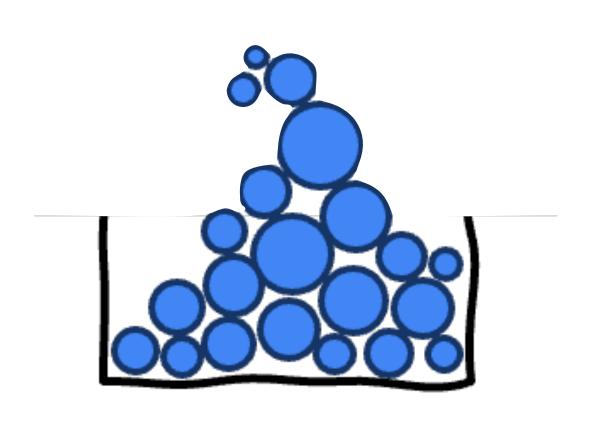
- Snow [Stomakhin et al. 2013]
- Phase-change [Tu et al. 2024] [Su et al. 2021] [Stomakhin et al. 2014]
- Granular
 [Li et al. 2022] [Tampubolon et al. 2017]
 [Klár et al. 2016]
 [Daviet and Bertails-Descoubes 2016]
- Magnetic [Sun et al. 2021]





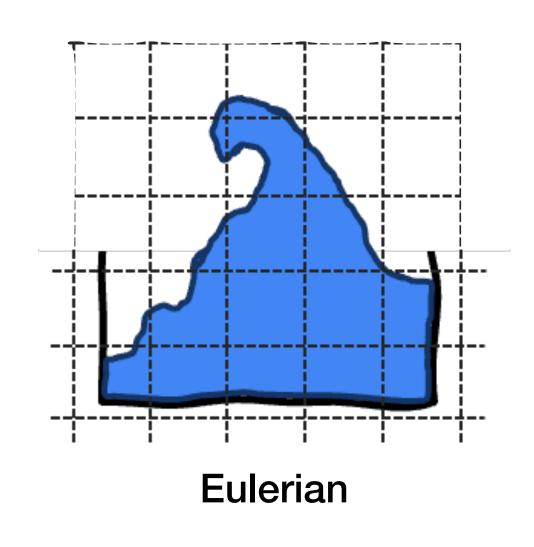


Part III - Summary



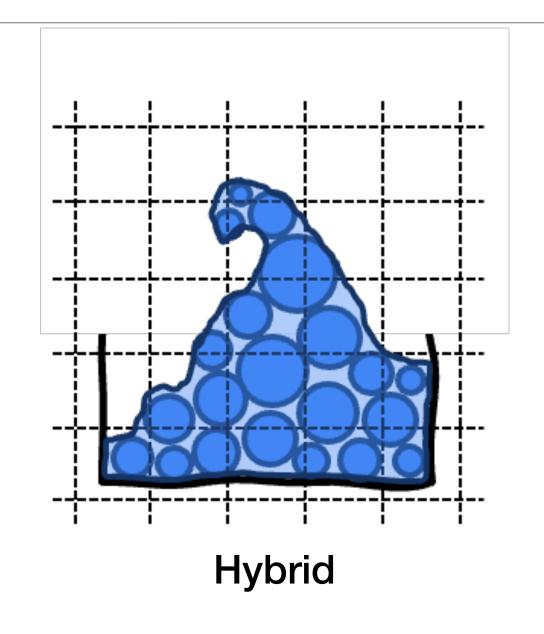
Lagrangian

- Topology changes
- **Y** Contact
- TFluid-rigid coupling





- **Turbulence**
- **Y** Bounded domains



- Topology changes
- Tomplex material models
- T Cutting and fracture





